Exact Travelling Wave Solutions for Stochastic Fractional Hirota-Satsuma Coupled KdV Equations

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Abstract. A modified fractional sub-equation method is used for constructing exact travelling wave solutions of nonlinear stochastic fractional partial differential equations. The main idea of this method is to take full advantage of the fractional Riccati equation, which has many exact solutions. Moreover, white noise functional solutions are obtained for the Wick-type stochastic fractional Hirota-Satsuma coupled KdV equations via Hermite transform and white noise analysis. These solutions include stochastic exponential decay, soliton and periodic wave solutions.

Keywords: Hirota-Satsuma coupled KdV equations; Fractional calculus; Riccati equation; White noise; Hermite transform.

PACS No.: 05.40. ± a, 02.30.Jr.

1 Introduction

In this paper, we consider the fractional Hirota-Satsuma coupled KdV equations given by

\[ \begin{aligned}
D^\alpha_t u &= p(t)D^{3\alpha}_x u + q(t)u D^\alpha_x u + r(t) D^\alpha_x (vw), \\
D^\alpha_t v &= -D^{3\alpha}_x v + 3uD^\alpha_x v, \\
D^\alpha_t w &= -D^{3\alpha}_x w + 3uD^\alpha_x w
\end{aligned} \]  

(1.1)
where \((x, t) \in \mathbb{R} \times \mathbb{R}^+_0, 0 < \alpha \leq 1\), \(D^\alpha_t u\) and \(D^\alpha_x u\) are the modified Riemann-Liouville derivatives and \(p(t), q(t)\) and \(r(t)\) are bounded measurable or integrable functions on \(\mathbb{R}^+_0\). As pointed in \([8, 21, 31]\) the interaction of two long waves with different dispersion relations can be described by Eq.(1.1). Moreover, if the problem is considered in random environment, we can get random fractional Hirota-Satsuma coupled KdV equations. In order to give the exact solutions of this random model, we only consider it in white noise environment, that is, we will study the following Wick-type stochastic fractional Hirota-Satsuma coupled KdV equations

\[
\begin{align*}
D^\alpha_t U &= P(t) \circ D^3_2 \alpha U + Q(t) \circ U \circ D^\alpha_t U + R(t) \circ D^\alpha_t (V \circ W), \\
D^\alpha_t V &= -D^3_2 \alpha V + 3U \circ D^\alpha_x V, \\
D^\alpha_t W &= -D^3_2 \alpha W + 3U \circ D^\alpha_x (V \circ W)
\end{align*}
\]  

(1.2)

where “\(\circ\)” is the Wick product on the Kondratiev distribution space \((S)_{-1}\) and \(P(t), Q(t)\) and \(R(t)\) are \((S)_{-1}\)-valued functions [22].

Many important phenomena in electromagnetics, acoustics, viscoelasticity, electrochemistry, cosmology and material science can be better described by fractional partial differential equations FPDEs [26, 28, 29]. Consequently, considerable attention has been given to the solution of the FPDEs. There are many methods for calculating the approximate solutions for nonlinear FPDEs such as the variational iterations method [32], Adomian decomposition method [2, 3], the homotopy perturbation method [17, 18] and the Exp-function method [25, 34, 35, 36, 39]. The exact solutions for nonlinear FPDEs are still under study until now. Li and He [24] introduced complex transform for reducing FPDEs into ordinary differential equations, so that all analytical methods for advanced calculus can be easily applied to fractional calculus. Also, He [19] introduced a new method to look for exact travelling wave solutions of nonlinear FPDEs. This method is called fractional sub-equation method, and it is based on the homogeneous balance principle [38] and Jumarie’s Riemann-Liouville derivatives [23].

It is well known that the solitons are stable against mutual collisions and behave like particles. In this sense, it is very important to study the nonlinear equations in random environment. However, variable coefficients nonlinear equations, as well as constant coefficients equations, cannot describe the realistic physical phenomena exactly. Wadati [33] first answered the interesting question, “How does external noise affect the motion of solitons?” and studied the diffusion of soliton of the KdV equation under Gaussian noise, which satisfies a diffusion equation in transformed coordinates. The cauchy problems associated with stochastic PDEs was discussed by many authors, e.g., de Bouard and Debussche [4, 5], Debussche and Printems [6, 7], Printems [30] and Ghany and Hyder [12]. On the basis of white noise functional analysis [22], Ghany et al. [9-16] studied more intensely the white noise functional solutions for some nonlinear stochastic PDEs.

The motivation of this paper is to propose a modified fractional sub-equation method to seek new exact travelling wave solutions for the variable coefficients fractional Hirota-Satsuma coupled KdV equations with the modified Riemann-Liouville derivatives . These solutions include exponential decay, soliton and periodic travelling wave solutions. Then, with the help of Hermit transform and white noise analysis, we employ these solutions to find new white noise functional solutions for the stochastic fractional Hirota-Satsuma coupled KdV equations in white noise environment, especially in its Wick version.

The rest of this paper is organized as follows: In Section 2, we recall the definition and some properties of the modified Riemann-Liouville derivative and introduce a modified fractional sube-
question method. In Section 3, we apply the modified fractional sub-equation method to explore exact travelling wave solutions for Eq.(1.1). In Section 4, we use the Hermite transform and [22,Theorem 4.1.1] to obtain white noise functional solutions for Eq.(1.2). In Section 5, we give illustrative examples for the investigated model. The last section is devoted to summary and discussion.

2 Preliminaries

There are different definitions for fractional derivatives, for more details see [28]. In our paper we use the modified Riemann-Liouville derivative defined by Jumarie [23]

\[
D_\alpha^x f(x) = \begin{cases} 
\frac{1}{\Gamma(1-\alpha)} \int_0^x (x-y)^{-\alpha-1}[f(y)-f(0)]dy, & \alpha < 0, \\
\frac{1}{\Gamma(1-\alpha)} \frac{d}{dx} \int_0^x (x-y)^{-\alpha}[f(y)-f(0)]dy, & 0 < \alpha < 1, \\
[f^{(\alpha-n)}(x)](n), & n \leq \alpha < n+1, \ n \in \mathbb{N}
\end{cases}
\]

(2.1)

which has merits over the original one, for example, the \(\alpha\)-order derivative of a constant is zero. Some properties of the modified Riemann-Liouville derivative were summarized in [23], three useful formulas of them are

\[
\begin{align*}
D_\alpha^x x^\beta &= \frac{\Gamma(1+\beta)}{\Gamma(1+\beta-\alpha)} x^{\beta-\alpha}, \quad \beta > 0, \\
D_\alpha^x (u(x)v(x)) &= u(x)D_\alpha^x v(x) + v(x)D_\alpha^x u(x), \\
D_\alpha^x [f(u(x))] &= \frac{df}{du} D_\alpha^x u(x) = \left(\frac{du}{dx}\right)^\alpha D^\alpha_\xi f(u).
\end{align*}
\]

(2.2)

Now, we outline the main idea of the modified fractional sub-equation method. Many authors considered nonlinear FPDE, say, in two variables

\[
F(u, u_x, u_t, D_\alpha^u u, D_\alpha^2 u, ...) = 0, \quad 0 < \alpha \leq 1
\]

(2.3)

where \(F\) is a nonlinear function with respect to the indicated variables. To determine the solution \(u = u(x, t)\) explicitly, we first introduce the following transformation

\[
u = u(\xi), \quad \xi = \xi(x, t)
\]

(2.4)

which converts Eq.(2.3) into a fractional ordinary differential equation

\[
G(u, u', u'', D_{\xi}^\alpha u, D_{\xi}^2 u, ...) = 0.
\]

(2.5)

Next we introduce a new variable \(Y = Y(\xi)\) which is a solution of the fractional Riccati equation

\[
D_\alpha^2 Y = h_0 + h_1 Y + h_2 Y^2, \quad 0 < \alpha \leq 1,
\]

(2.6)

where \(h_0, h_1\) and \(h_2\) are arbitrary constants. Eq.(2.6) is the fractional Riccati differential equation, where \(\alpha\) is a parameter describing the order of the fractional derivative. In the case of \(\alpha = 1\) Eq.(2.6)
is reduced to the classical Riccati differential equation. The importance of this equation usually arises in the optimal control problems. The feed back gain of the linear quadratic optimal control depends on a solution of a Riccati differential equation which has to be found for the whole time horizon of the control process [27, 38]. Then we propose the following series expansion as a solution of Eq.(2.3)

\[ u(x, t) = u(\xi) = \sum_{k=0}^{n} a_k(x, t)Y^k(\xi) + \sum_{k=1}^{n} b_k(x, t)Y^{-k}(\xi), \quad (2.7) \]

where \(a_k(k = 0, 1, ..., n), b_k(k = 1, ..., n)\) are functions to be determined later and \(n\) is a positive integer which can be determined via the balancing of the highest derivative term with the nonlinear term in equation Eq.(2.5). Inserting Eq.(2.7) into Eq.(2.5) and using Eq.(2.6) will give an algebraic equation in powers of \(Y\). Since all coefficients of \(Y^k\) must vanish, this will give a system of algebraic equations with respect to \(a_k\) and \(b_k\). With the aid of Mathematica, we can determine \(a_k\) and \(b_k\).

According to the recent paper by Zhang et al. [38], we can deduce the following set of solutions of Eq.(2.3) according to the recent paper by Zhang et al. [38], we can deduce the following set of solutions of Eq.(2.3)

\[
\begin{align*}
Y_1(\xi) &= E_\alpha(\xi) - 1, \\
Y_2(\xi) &= \coth_\alpha(\xi) \pm \csch_\alpha(\xi), \\
Y_3(\xi) &= \tanh_\alpha(\xi) \pm i \sech_\alpha(\xi), \\
Y_4(\xi) &= \frac{1}{2} \tan_\alpha(2\xi), \\
Y_5(\xi) &= \frac{1}{2} \cot_\alpha(2\xi),
\end{align*}
\]

(2.8)

with the generalized hyperbolic and trigonometric functions

\[
\begin{align*}
\tanh_\alpha(x) &= \frac{\sinh_\alpha(x)}{\cosh_\alpha(x)}, & \coth_\alpha(x) &= \frac{\cosh_\alpha(x)}{\sinh_\alpha(x)}, & \csch_\alpha(x) &= \frac{1}{\sinh_\alpha(x)}, & \sech_\alpha(x) &= \frac{1}{\cosh_\alpha(x)}, \\
\sinh_\alpha(x) &= \frac{E_\alpha(x^\alpha) - E_\alpha(-x^\alpha)}{2}, & \cosh_\alpha(x) &= \frac{E_\alpha(x^\alpha) + E_\alpha(-x^\alpha)}{2}, & \tan_\alpha(x) &= \frac{\sin_\alpha(x)}{\cos_\alpha(x)}, \\
\cot_\alpha(x) &= \frac{\cos_\alpha(x)}{\sin_\alpha(x)}, & \sin_\alpha(x) &= \frac{E_\alpha(ix^\alpha) - E_\alpha(-ix^\alpha)}{2i}, & \cos_\alpha(x) &= \frac{E_\alpha(ix^\alpha) + E_\alpha(-ix^\alpha)}{2},
\end{align*}
\]

defined by the Mittag-Leffler function \(E_\alpha(y) = \sum_{j=0}^{\infty} \frac{y^j}{\Gamma(1 + j\alpha)}\). For more details about the generalized exponential, hyperbolic and trigonometric functions see [29].

3 Exact Travelling Wave Solutions for Eq.(1.1)

In this section, we apply Hermite transform, white noise analysis, and modified fractional sub-equation method to explore exact travelling wave solutions for Eq.(1.1). Taking the Hermite transform of Eq.(1.2), we get the deterministic system

\[
\begin{align*}
D_t^\alpha \widetilde{U}(x, t, z) &= \widetilde{P}(t, z)D_x^{3\alpha} \widetilde{U}(x, t, z) + \widetilde{Q}(t, z) \widetilde{U}(x, t, z) D_x^2 \widetilde{U}(x, t, z), \\
+\widetilde{R}(t, z) D_x^\alpha (\widetilde{V}(x, t, z) \widetilde{W}(x, t, z)), \\
D_t^\alpha \widetilde{V}(x, t, z) &= -D_x^{3\alpha} \widetilde{V}(x, t, z) + 3 \widetilde{U}(x, t, z) D_x^\alpha \widetilde{V}(x, t, z), \\
D_t^\alpha \widetilde{W}(x, t, z) &= -D_x^{3\alpha} \widetilde{W}(x, t, z) + 3 \widetilde{U}(x, t, z) D_x^\alpha \widetilde{W}(x, t, z),
\end{align*}
\]

(3.1)
where $z = (z_1, z_2, \ldots) \in (\mathbb{C}^N)_c$ is a vector parameter. To look for the travelling wave solution of Eq.(3.1), we make the transformations $u(x, t, z) := \tilde{U}(x, t, z) = \phi (\xi(x, t, z))$, $v(x, t, z) := \tilde{V}(x, t, z) = \psi (\xi(x, t, z))$ and $w(x, t, z) := \tilde{W}(x, t, z) = \chi (\xi(x, t, z))$, with

$$\xi(x, t, z) = kx + s \int_0^t \lambda(\tau, z) d\tau + c,$$

where $k, s$ and $c$ are arbitrary constants which satisfy $ks \neq 0$, $\lambda(t, z)$ is a non zero functions of the indicated variables to be determined. So, Eq.(3.1) can be changing into the form

$$\begin{align*}
(sl)^\alpha D_\xi^\alpha \phi &= k^{3\alpha} p D_x^{2\alpha} \phi + k^\alpha q \phi D_\xi^\alpha \phi + k^\alpha r D_\xi^\alpha (\psi \chi), \\
(sl)^\alpha D_\xi^\alpha \psi &= -k^{3\alpha} D_x^{3\alpha} \psi + 3k^\alpha \phi D_\xi^\alpha \psi, \\
(sl)^\alpha D_\xi^\alpha \chi &= -k^{3\alpha} D_x^{3\alpha} \chi + 3k^\alpha \phi D_\xi^\alpha \chi,
\end{align*}$$

(3.2)

where $p(t, z) := \bar{P}(t, z)$, $q(t, z) := \bar{Q}(t, z)$ and $r(t, z) := \bar{R}(t, z)$. Balancing the highest order linear terms and nonlinear terms in Eq.(3.2), gives the following ansätze

$$\begin{align*}
u(x, t, z) &= a_0(t, z) + a_1(t, z) Y(\xi) + a_2(t, z) Y^2(\xi) + b_1(t, z) Y^{-1}(\xi) + b_2(t, z) Y^{-2}(\xi), \\
v(x, t, z) &= c_0(t, z) + c_1(t, z) Y(\xi) + c_2(t, z) Y^2(\xi) + d_1(t, z) Y^{-1}(\xi) + d_2(t, z) Y^{-2}(\xi), \\
w(x, t, z) &= e_0(t, z) + e_1(t, z) Y(\xi) + e_2(t, z) Y^2(\xi) + f_1(t, z) Y^{-1}(\xi) + f_2(t, z) Y^{-2}(\xi)
\end{align*}$$

(3.3)

where $Y(\xi)$ satisfies the fractional Riccati equation (2.6). By substituting Eq.(3.3) along with Eq.(2.6) into Eq.(3.2), collect the coefficients of $Y^k (k = -5, -4, \ldots, 5)$ and set them to be zero, we will obtain a system of algebraic equations in the unknowns $a_k, c_k, e_k (k = 0, 1, 2)$, $b_k, d_k, f_k (k = 1, 2)$ and $l$ of the form

$$\begin{align*}
(sl)^\alpha G_i^\phi &= k^{3\alpha} p K_i^\phi + k^\alpha q \zeta_i^\phi + k^\alpha r (\rho_i^1 + \rho_i^2), & i &= 0, 1, 2, 3, \\
(sl)^\alpha H_i^\phi &= k^{3\alpha} p L_i^\phi + k^\alpha q \eta_i^\phi + k^\alpha r (\lambda_i^1 + \lambda_i^2), & i &= 1, 2, 3, \\
k^{2\alpha} p K_i^\phi + q \zeta_i^\phi + r (\rho_i^1 + \rho_i^2) &= 0, & i &= 4, 5, \\
k^{2\alpha} p L_i^\phi + q \eta_i^\phi + r (\lambda_i^1 + \lambda_i^2) &= 0, & i &= 4, 5, \\
(sl)^\alpha G_i^\psi &= -k^{3\alpha} K_i^\psi + 3k^\alpha \zeta_i^\psi, & i &= 0, 1, 2, 3, \\
(sl)^\alpha H_i^\psi &= -k^{3\alpha} L_i^\psi + 3k^\alpha \eta_i^\psi, & i &= 1, 2, 3, \\
-k^{2\alpha} K_i^\psi + 3 \zeta_i^\psi &= 0, & i &= 4, 5, \\
-k^{2\alpha} L_i^\psi + 3 \eta_i^\psi &= 0, & i &= 4, 5, \\
(sl)^\alpha G_i^\chi &= -k^{3\alpha} K_i^\chi + 3k^\alpha \zeta_i^\chi, & i &= 0, 1, 2, 3, \\
(sl)^\alpha H_i^\chi &= -k^{3\alpha} L_i^\chi + 3k^\alpha \eta_i^\chi, & i &= 1, 2, 3, \\
-k^{2\alpha} K_i^\chi + 3 \zeta_i^\chi &= 0, & i &= 4, 5, \\
-k^{2\alpha} L_i^\chi + 3 \eta_i^\chi &= 0, & i &= 4, 5.
\end{align*}$$

(3.4)

where, $G_0^\phi = h_0 a_1 - h_2 b_1$, $G_1^\phi = 2h_0 a_2 + h_1 a_1$, $G_2^\phi = 2h_1 a_2 + h_2 a_1$, $G_3^\phi = 2h_2 a_2$, $G_0^\psi = h_0 c_1 - h_2 d_1$, $G_1^\psi = 2h_0 c_2 + h_1 c_1$, $G_2^\psi = 2h_1 c_2 + h_2 c_1$, $G_3^\psi = 2h_2 c_2$, $G_0^\chi = h_0 e_1 - h_2 f_1$, $G_1^\chi = 2h_0 e_2 + h_1 e_1$, $G_2^\chi = 2h_1 e_2 + h_2 e_1$, $G_3^\chi = 2h_2 e_2$, $H_1^\phi = -(2h_2 b_2 + h_1 b_1)$, $H_2^\phi = -(2h_1 b_2 +
In the remaining part of this section we investigate and solve our problem for some particular cases for the fractional Riccati equation (2.6) as follows.

**Case A.** If we set \( h_0 = h_1 = 1, \) \( h_2 = 0 \) in Eq. (2.6), and use Mathematica to solve the result-
ing system, we will obtain the following set of solutions

**Set 1:** \( a_1 = a_2 = c_1 = c_2 = e_1 = e_2 = 0, \ a_0 = \frac{pk^{2\alpha}}{q} - \frac{4rp}{9q} \left( \frac{6r}{p} - \frac{21p}{q} \right)^{-1} k^{2\alpha}, \ c_0 = e_0 = \frac{pk^{2\alpha}}{3}, \ b_1 = \left( \frac{6r}{p} - \frac{21p}{q} \right) k^{2\alpha}, \ b_2 = -\frac{6pk^{2\alpha}}{q}, \ d_1 = f_1 = \frac{2}{3} k^{2\alpha}, \ d_2 = f_2 = 4k^{2\alpha}, \ l = \left( \frac{2pk^{3\alpha}}{s^{\alpha}} \right)^{\frac{1}{\pi}} \), under condition that \( 77p^2 - 64qr = 0 \).

Substituting these values in Eq.(3.3) and using Eq.(2.8), we obtain the following exponential decay wave solution of Eq.(3.1):

\[
\begin{align*}
  u_1(x, t, z) &= \frac{k^{2\alpha} p(t, z)}{q(t, z)} - \frac{4k^{2\alpha} r(t, z) p^2(t, z)}{189q(t, z) q(t, z) - 54r(t, z) q(t, z) - 21p^2(t, z)} \\
  &\quad + \frac{k^{2\alpha} \left[ (6r(t, z) q(t, z) - 21p^2(t, z)) \left( Y_1[\xi_1(x, t, z)] \right) \right] - 6p^2(t, z)}{p(t, z) q(t, z) \left( Y_1[\xi_1(x, t, z)] \right)^2}, \tag{3.5}
\end{align*}
\]

\[
\begin{align*}
  v_1(x, t, z) &= \frac{k^{2\alpha}}{3} \ p(t, z) + \frac{2k^{2\alpha} \left( Y_1[\xi_1(x, t, z)] + 6 \right)}{3 \left( Y_1[\xi_1(x, t, z)] \right)^2}, \tag{3.6}
\end{align*}
\]

\[
\begin{align*}
  w_1(x, t, z) &= \frac{k^{2\alpha}}{3} \ p(t, z) + \frac{2k^{2\alpha} \left( Y_1[\xi_1(x, t, z)] + 6 \right)}{3 \left( Y_1[\xi_1(x, t, z)] \right)^2}, \tag{3.7}
\end{align*}
\]

with

\[
\xi_1(x, t, z) = kx + (2k^{3\alpha})^{\frac{1}{\pi}} \int_0^t \left( p(\tau, z) \right)^{\frac{1}{\pi}} d\tau + c_1.
\]

**Case B.** If we set \( h_0 = -h_2 = \frac{1}{2}, \ h_1 = 0 \) in Eq.(2.6), and use Mathematica to solve the resulting system, we will obtain the following sets of solutions

**Set 2:** \( a_1 = b_1 = c_1 = c_2 = d_1 = e_1 = e_2 = f_1 = 0, \ a_0 = \frac{(2 + 3p)k^{2\alpha}}{q}, \ c_0 = -e_0 = \frac{pk^{2\alpha}}{3}, \ a_2 = 2b_2 = \frac{-3pk^{2\alpha}}{q}, \ f_2 = d_2 = k^{2\alpha}, \ l = \left( \frac{(2 + p)k^{3\alpha}}{s^{\alpha}} \right)^{\frac{1}{\pi}} \).

Substituting these values in Eq.(3.3) and using Eq.(2.8), we obtain the following soliton wave solutions of Eq.(3.1)

\[
\begin{align*}
  u_i(x, t, z) &= \frac{(2 + 3p(t, z))k^{2\alpha}}{q(t, z)} - \frac{3p(t, z)k^{2\alpha}}{2q(t, z)} \left( 2Y_i^2[\xi_2(x, t, z)] + Y_i^{-2}[\xi_2(x, t, z)] \right), \tag{3.8}
\end{align*}
\]

\[
\begin{align*}
  v_i(x, t, z) &= \frac{p(t, z)k^{2\alpha}}{3} + k^{2\alpha} Y_i^{-2}[\xi_2(x, t, z)], \tag{3.9}
\end{align*}
\]

\[
\begin{align*}
  w_i(x, t, z) &= \frac{p(t, z)k^{2\alpha}}{3} + k^{2\alpha} Y_i^{-2}[\xi_2(x, t, z)], \tag{3.10}
\end{align*}
\]

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with $i = 2, 3$ and
\[
\xi_2(x, t, z) = kx + k^3 \int_0^t (2 + p(\tau, z)) \frac{1}{3} \, d\tau + c_2.
\]

**Set 3:** $a_1 = a_2 = b_1 = c_1 = c_2 = d_1 = e_1 = e_2 = f_1 = 0$, $a_0 = \frac{(2 + 4p)k^{2\alpha}}{q}$, $c_0 = -e_0 = \frac{pk^{2\alpha}}{3}$, $b_2 = \frac{-3pk^{2\alpha}}{2q}$, $f_2 = k^2$, $l = \left(\frac{(2 + p)k^{2\alpha}}{s^\alpha}\right)^{\frac{1}{\alpha}}$.

Substituting these values in Eq.(3.3) and using Eq.(2.8), we obtain the following soliton wave solutions of Eq.(3.1)
\[
u_{i+2}(x, t, z) = \frac{(2 + 4p(t, z))k^{2\alpha}}{q(t, z)} - \frac{3p(t, z)k^{2\alpha}}{2q(t, z)}Y_i^{-2}[\xi_2(x, t, z)],
\]
(3.11)
\[
v_{i+2}(x, t, z) = \frac{p(t, z)k^{2\alpha}}{3} + k^2Y_i^{-2}[\xi_2(x, t, z)],
\]
(3.12)
\[
w_{i+2}(x, t, z) = \frac{p(t, z)k^{2\alpha}}{3} + k^2Y_i^{-2}[\xi_2(x, t, z)]
\]
(3.13)

with $i = 2, 3$.

**Case C.** If we set $h_0 = \frac{1}{4}h_2 = 1$, $h_1 = 0$ in Eq.(2.6), and use Mathematica to solve the resulting system, we will obtain the following set of solutions

**Set 4:** $a_1 = b_1 = c_1 = d_1 = d_2 = e_1 = f_2 = 0$, $a_0 = \frac{(32 + 31p)k^{2\alpha}}{q} - \frac{4rk^{2\alpha}}{9}$, $a_2 = -\frac{96pk^{2\alpha}}{q}$, $e_0 = \frac{pk^{2\alpha}}{3}$, $c_2 = e_2 = 64k^{2\alpha}$, $l = \left(\frac{(p - 32)k^{2\alpha}}{s^\alpha}\right)^{\frac{1}{\alpha}}$.

Substituting these values in Eq.(3.3) and using Eq.(2.8), we obtain the following periodic wave solutions of Eq.(3.1)
\[
u_{i+2}(x, t, z) = \frac{(32 + 31p(t, z))k^{2\alpha}}{q(t, z)} - \frac{4r(t, z)k^{2\alpha}}{9} - \frac{96p(t, z)k^{2\alpha}}{q(t, z)}Y_i^{-2}[\xi_3(x, t, z)],
\]
(3.14)
\[
u_{i+2}(x, t, z) = \frac{p(t, z)k^{2\alpha}}{3} + 64k^{2\alpha}Y_i^{-2}[\xi_3(x, t, z)],
\]
(3.15)
\[
w_{i+2}(x, t, z) = \frac{p(t, z)k^{2\alpha}}{3} + 64k^2Y_i^{-2}[\xi_3(x, t, z)]
\]
(3.16)
with \( i = 4, 5 \) and
\[
\xi_3(x, t, z) = kx + k^3 \int_0^t (p(\tau, z) - 32)^{\frac{1}{2}} d\tau + c_3.
\]

At the end of this section we should remark that, there exists an infinitely number of exact travelling wave solutions for Eq.(1.1); these solutions come from solving the system (3.4) with regard to the fractional Riccati equation (2.6). The above mentioned cases are just to clarify how far our technique is applicable.

## 4 White Noise Functional Solutions of Eq.(1.2)

In this section, we use the inverse Hermite transform and [22,Theorem 4.1.1] to obtain white noise functional solutions for Eq.(1.2). The properties of generalized exponential, hyperbolic and trigonometric functions yield that there exists a bounded open set \( G \subset \mathbb{R} \times \mathbb{R}_+ \), \( m < \infty \), \( n > 0 \) such that the solution \( \{u(x, t, z), v(x, t, z), w(x, t, z)\} \) of Eq.(3.1) and all its fractional derivatives which are involved in Eq.(3.1) are uniformly bounded for \( (x, t, z) \in G \times K_m(n) \), continuous with respect to \( (x, t) \in G \) for all \( z \in K_m(n) \) and analytic with respect to \( z \in K_m(n) \), for all \( (x, t) \in G \). From [22,Theorem 4.1.1], there exist \( U(x, t), V(x, t), W(x, t) \in (S)_-1 \) such that \( u(x, t, z) = \tilde{U}(x, t)(z) \), \( v(x, t, z) = \tilde{V}(x, t)(z) \) and \( w(x, t, z) = \tilde{W}(x, t)(z) \) for all \( (x, t, z) \in G \times K_m(n) \) and \( \{U(x, t), V(x, t), W(x, t)\} \) solves (in the strong sense in \((S)_-1\) Eq.(1.2) in \((S)_-1\). Hence, for \( P(t)Q(t)R(t) \neq 0 \), the white noise functional solutions of Eq.(1.2) can be obtained by applying the inverse Hermite transform to Eqs.(3.5)-(3.16) as follows

- **Stochastic exponential decay wave solution**:
  \[
  U_1(x, t) = \frac{k^{2\alpha} P(t)}{Q(t)} - \frac{4k^{2\alpha} R(t) \circ P^{\circ 2}(t)}{54R(t) \circ Q(t) - 189P^{\circ 2}(t)} + \frac{k^{2\alpha} \left[ (6R(t) \circ Q(t) - 21P^{2}(t)) \circ \left( Y_1^{\circ} \right) \right]}{P(t) \circ Q(t) \circ \left( Y_1^{\circ} \right)}^{\circ 2},
  \]  
  \( (4.1) \)

- **Stochastic soliton wave solutions**:
  \[
  V_1(x, t) = \frac{k^{2\alpha}}{3} P(t) + \frac{2k^{2\alpha} Y_1^{\circ} \left[ \Xi_1(x, t) \right]}{3(Y_1^{\circ} \left[ \Xi_1(x, t) \right])^{\circ 2}},
  \]  
  \( (4.2) \)

  \[
  W_1(x, t) = \frac{k^{2\alpha}}{3} P(t) + \frac{2k^{2\alpha} (Y_1^{\circ} \left[ \Xi_1(x, t) \right] + 6)}{3(Y_1^{\circ} \left[ \Xi_1(x, t) \right])^{\circ 2}},
  \]  
  \( (4.3) \)

with
\[
\Xi_1(x, t, z) = kx + (2k^{3\alpha})^{\frac{1}{\alpha}} \int_0^t (P(\tau))^\frac{1}{\alpha} d\tau + c_1.
\]

- **Stochastic soliton wave solutions**:
  \[
  U_i(x, t) = \frac{k^{2\alpha} (2 + 3P(t))}{Q(t, z)} - \frac{3k^{2\alpha} P(t)}{2Q(t)} \circ (2Y_i^{\circ 2} \left[ \Xi_2(x, t) \right] + Y_i^{\circ(-2)} \left[ \Xi_2(x, t) \right]),
  \]  
  \( (4.4) \)
\[ V_i(x,t) = \frac{k^{2\alpha}P(t)}{3} + k^{2\alpha}Y_i^{\alpha(-2)}[\Xi_2(x,t)], \quad (4.5) \]

\[ W_i(x,t) = \frac{k^{2\alpha}P(t)}{3} + k^{2\alpha}Y_i^{\alpha(-2)}[\Xi_2(x,t)], \quad (4.6) \]

\[ U_{i+2}(x,t) = \frac{k^{2\alpha}(2 + 4P(t))}{Q(t)} - \frac{3k^{2\alpha}P(t)}{2Q(t)} \cdot Y_i^{\alpha(-2)}[\Xi_2(x,t)], \quad (4.7) \]

\[ V_{i+2}(x,t) = \frac{k^{2\alpha}P(t)}{3} + k^{2\alpha}Y_i^{\alpha(-2)}[\Xi_2(x,t)], \quad (4.8) \]

\[ W_{i+2}(x,t) = \frac{k^{2\alpha}P(t)}{3} + k^{2\alpha}Y_i^{\alpha(-2)}[\Xi_2(x,t)], \quad (4.9) \]

with \( i = 2, 3 \) and

\[ \Xi_2(x,t) = kx + k^3 \int_0^t (2 + P(\tau))^{\frac{3}{2}} d\tau + c_2. \]

- Stochastic periodic wave solutions:

\[ U_{i+2}(x,t) = \frac{k^{2\alpha}(32 + 31P(t))}{Q(t)} - \frac{4k^{2\alpha}R(t)}{9} - \frac{96k^{2\alpha}P(t)}{Q(t)} \cdot Y_i^{\alpha2}[\Xi_3(x,t)], \quad (4.10) \]

\[ V_{i+2}(x,t) = \frac{k^{2\alpha}P(t)}{3} + 64k^{2\alpha}Y_i^{\alpha(-2)}[\Xi_3(x,t)], \quad (4.11) \]

\[ W_{i+2}(x,t) = \frac{k^{2\alpha}P(t)}{3} + 64k^{2\alpha}Y_i^{\alpha(-2)}[\Xi_3(x,t)], \quad (4.12) \]

with \( i = 4, 5 \) and

\[ \Xi_3(x,t) = kx + k^3 \int_0^t (P(\tau) - 32)^{\frac{3}{2}} d\tau + c_3. \]

We observe that for different forms of \( P(t), Q(t) \) and \( R(t) \), we can get different solutions of Eq.(1.2) from Eqs.(4.1)-(3.12).

5 Examples

Let \( \tilde{B}_t \) be the Gaussian white noise, where \( B_t \) is Brownian motion. We have the Hermite transform \( \tilde{B}(t,z) = \sum_{k=1}^{\infty} z_k \int_0^t \pi_k(s) ds \). Since \( E_{\alpha}(B_t) = E_{\alpha}(B_t - t^2/2) \), we have \( \tan_{\alpha}(B_t) = \tan_{\alpha}(B_t - t^2/2) \), \( \cot_{\alpha}(B_t) = \cot_{\alpha}(B_t - t^2/2) \), \( \tanh_{\alpha}(B_t) = \tanh_{\alpha}(B_t - t^2/2) \), \( \coth_{\alpha}(B_t) = \coth_{\alpha}(B_t - t^2/2) \), \( \sech_{\alpha}(B_t) = \sech_{\alpha}(B_t - t^2/2) \) and \( \csch_{\alpha}(B_t) = \csch_{\alpha}(B_t - t^2/2) \). Suppose \( P(t) = R(t) = \theta_1Q(t) \) and \( Q(t) = (q(t) + \theta_2B_t)^{\alpha} \), where \( \theta_1, \theta_2 \) are arbitrary constants and \( q(t) \) is integrable or
bounded measurable function on $\mathbb{R}_+$. Hence, for $Q(t) \neq 0$ the white noise functional solutions of Eq.(1.2) in non-Wick version are as follows:

$$U_8(x,t) = \theta_1 k^{2\alpha} - \frac{4\theta_1 k^{2\alpha} Q(t)}{9(6 - 21\theta_1)} + k^{2\alpha} [(6 - 21\theta_1) Y_1[\Xi_1^*(x,t)] - 6\theta_1 Y_1^{-2}[\Xi_1^*(x,t)],$$  \hspace{1cm} (5.1) 

$$V_8(x,t) = \frac{\theta_1 k^{2\alpha} Q(t)}{3} + \frac{2k^{2\alpha}}{3} (Y_1[\Xi_1^*(x,t)] + 6) Y_1^{-2}[\Xi_1^*(x,t)],$$  \hspace{1cm} (5.2) 

$$W_8(x,t) = \frac{\theta_1 k^{2\alpha} Q(t)}{3} + \frac{2k^{2\alpha}}{3} (Y_1[\Xi_1^*(x,t)] + 6) Y_1^{-2}[\Xi_1^*(x,t)],$$  \hspace{1cm} (5.3) 

with

$$\Xi_1^*(x,t) = kx + (2k^3)^{1/\alpha} \theta_1 \left( \int_0^t q(\tau)d\tau - \theta_2 B_t + \frac{\theta_2 t^2}{2} \right) + c_1,$$

$$U_{i+7}(x,t) = \frac{k^{2\alpha}(2 + 3\theta_1 Q(t))}{Q(t)} - \frac{3\theta_1 k^{2\alpha}}{2} (2Y_i^2[\Xi_2^*(x,t)] + Y_i^{-2}[\Xi_2^*(x,t)]),$$  \hspace{1cm} (5.4) 

$$V_{i+7}(x,t) = \frac{\theta_1 k^{2\alpha} Q(t)}{3} + k^{2\alpha} Y_i^{-2}[\Xi_2^*(x,t)],$$  \hspace{1cm} (5.5) 

$$W_{i+7}(x,t) = \frac{\theta_1 k^{2\alpha} Q(t)}{3} + k^{2\alpha} Y_i^{-2}[\Xi_2^*(x,t)],$$  \hspace{1cm} (5.6) 

$$U_{i+9}(x,t) = \frac{k^{2\alpha}(2 + 4\theta_1 Q(t))}{Q(t)} + \frac{3\theta_1 k^{2\alpha}}{2} Y_i^{-2}[\Xi_2^*(x,t)],$$  \hspace{1cm} (5.7) 

$$V_{i+9}(x,t) = \frac{k^{2\alpha}\theta_1 Q(t)}{3} - k^{2\alpha} Y_i^{-2}[\Xi_2^*(x,t)],$$  \hspace{1cm} (5.8) 

$$W_{i+9}(x,t) = \frac{k^{2\alpha}\theta_1 Q(t)}{3} - k^{2\alpha} Y_i^{-2}[\Xi_2^*(x,t)],$$  \hspace{1cm} (5.9) 

with $i = 2, 3$ and

$$\Xi_2^*(x,t) = kx + k^3 \left( \int_0^t q(\tau)d\tau - \theta_2 B_t + \frac{\theta_2 t^2}{2} \right) + c_2,$$

$$U_{i+9}(x,t) = \frac{k^{2\alpha}(32 + 31\theta_1 Q(t))}{Q(t)} + \frac{4k^{2\alpha}\theta_1 Q(t)}{9} - 96\theta_1 k^{2\alpha} Y_i^{-2}[\Xi_2^*(x,t)],$$  \hspace{1cm} (5.10)
\[ V_{i+9}(x,t) = \frac{k^{2\alpha} \theta_t Q(t)}{3} - 64k^{2\alpha} Y_i^{-2}[\Xi_3^*(x,t)], \quad (5.11) \]

\[ V_{i+9}(x,t) = \frac{k^{2\alpha} \theta_t Q(t)}{3} - 64k^{2\alpha} Y_i^{-2}[\Xi_3^*(x,t)], \quad (5.12) \]

with \( i = 4, 5 \) and

\[ \Xi_3^*(x,t) = kx + k^3 \left( 30 \int_0^t q(\tau)d\tau + \theta_2 B_t - \frac{\theta_2 t^2}{2} \right) + c_3. \]

6 Summary and Discussion

Our first interest in this work is to implement new strategies that give white noise functional solutions of the variable coefficients Wick-type stochastic fractional Hirota-Satsuma coupled KdV equations. The strategies that will be pursued in this work rest mainly on Hermite transform, white noise theory and modified fractional sub-equation method, all of which are employed to find white noise functional solutions of Eq.(1.2). The proposed schemes, as we believe, are entirely new and introduce new solutions in addition to the well-known traditional solutions. The ease of using these methods are showed its power to determine shock or solitary type of solutions. Obviously, the planner which we have proposed in this paper can be also applied to other nonlinear PDEs in mathematical physics such as KdV-Burgers, modified KdV-Burgers, Sawada-Kotera, Zhiber-Shabat and Benjamin-Bona-Mahony equations. Note that, if \( \alpha = 1 \), Eq.(1.2) is reduced to the stochastic generalized Hirota-Satsuma coupled KdV equations. Also, if \( \alpha = 1 \), \( p = 0.5 \) and \( r = -q = 3 \), Eq.(1.1) is reduced to the generalized Hirota-Satsuma coupled KdV equations. Hence, our results can be considered a generalization of the work due to Ghany [9] and Wu [37]. Moreover, there is a unitary mapping between the Gaussian white noise space and the Poisson white noise space, this connection was given by Benth and Gjerde [1]. Hence, with the help of this connection, we can derive some Poisson white noise functional solutions, if the coefficients \( P(t), Q(t) \) and \( R(t) \) are Poisson white noise functions in Eq.(1.2).

References


